

So far: while the graphs were both discrete and continuous, but all the boundaries were continuous.

Today we study **quantizations**. Some meanings:

- Continuous set \rightarrow discrete approximation

- Deformation of weights $1 \rightarrow q^v$
quantization parameter \swarrow \nwarrow some function of path

- Deformation of objects: sets \rightarrow subspaces
groups \rightarrow quantum groups

We start from **q -Pascal graph**.

Origins:

- quasi-invariant measures on 0/1 sequences
- $GL(\infty, q)$ -invariant measures on subspaces

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A q -Analogue of de Finetti's Theorem

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 PDF

Definition: Fix $q > 0$. We say that a random sequence $(z_i)_{i \geq 1} \in \{0,1\}^{\infty}$ is q -exchangeable if

$$\text{Prob}(z_1 = x_1, \dots, z_{i-1} = x_{i-1}, \underbrace{z_i = x_{i+1}, z_{i+1} = x_i}_{\text{swapped}}, z_{i+2} = x_{i+2}, \dots, z_n = x_n) \\ = q^{x_i - x_{i+1}} \text{Prob}(z_1 = x_1, \dots, z_n = x_n)$$

[For all $1 \leq i < n$ and all $x_1, \dots, x_n \in \{0,1\}$]

Meaning: At $q = 1$, the measure is $S(\infty)$ -invariant. At general $q > 0$, it is quasi-invariant: changes by multiplication with a controlled constant.

Example: Suppose $\text{Prob}(1, 0, 0) = c$, then $\text{Prob}(0, 1, 0) = c \cdot q$ and $\text{Prob}(0, 0, 1) = c \cdot q^2$. Moving 1 to the right multiplies the probability by q .

Exercise 1: As usual, we can identify 0/1 sequences with paths in the Pascal graph. At $q=1$, invariant measure corresponds to a central measure on paths [All paths with the same end-point have the same probability]. At general $q > 0$ the notion of centrality changes: now given the end-point, the probability of a path is proportional to the product of weights over its edges. Find these weights.

Observation: If we swap $(z_1, z_2, \dots) \leftrightarrow (1-z_1, 1-z_2, \dots)$ and simultaneously invert $q \leftrightarrow q^{-1}$, the notion of q -exchangeability remains the same.

Hence, we can assume without loss of generality that

$$0 < q < 1$$

Theorem 1: Assume $0 < q < 1$. Extreme q -exchangeable probability measures on $0/1$ sequences are parameterized by points of $\{0, 1, 2, \dots\} \cup \{\infty\}$. They are:

- ∞ -measure assigns probability 1 to $(1, 1, 1, \dots)$
- 0 -measure assigns probability 1 to $(0, 0, 0, \dots)$
- k -measure is supported on sequences with k ones at positions $0 < y_1 < y_2 < \dots < y_k$

Denoting also $y_0 = 0$, the spacings $y_{i+1} - y_i - 1$ are independent and geometrically distributed:

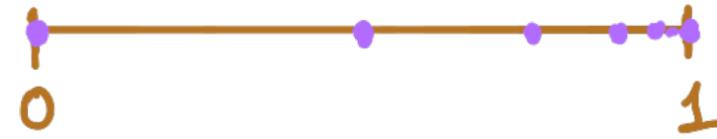
$$\text{Prob}(y_{i+1} - y_i - 1 = m) = (1 - q^{k-i}) \cdot q^{m \cdot (k-i)}, \quad \begin{array}{l} m = 0, 1, 2, \dots \\ 0 \leq i \leq k-1 \end{array}$$

Sharp contrast with $q = 1$:

- At $q = 1$ all but two measures have infinitely many 0s and 1s
- At $q \neq 1$ all but one measure have finitely many 1s
.... However!

Theorem 2: Identify measures of Theorem 1 with points

$$\{1\} \cup \{1 - q^k\}_{k \geq 0} \subset [0, 1]$$



Then as $q \rightarrow 1$ the points fill entire segment $[0, 1]$ and the measures converge to i.i.d. Bernoulli measures of parameter $p \in [0, 1]$ ($\text{Prob}(Z_i = 1) = p$)

Proof of Theorem 2: Assume that $q \rightarrow 1$ and $k \rightarrow \infty$ in such a way that $q^k \rightarrow 1 - p$. We need to show that the measures of Theorem 1 converge to i.i.d. p -Bernoulli.

Indeed, the spacing between 1's becomes $\text{Geom}(q^k) \rightarrow \text{Geom}(1-p)$

But $\text{Geom}(1-p)$ is precisely the spacing between 1's in i.i.d. p -Bernoulli: $\text{Prob}(\underbrace{0 \dots 0}_m 1) = p \cdot (1-p)^m$, $m \geq 0$.



Sketch of the proof of Theorem 1: The supports of the given measures are disjoint, hence, it suffices to show that they form the Martin boundary.

Thus, we need to choose $0 \leq k(n) \leq n$ for $n=1, 2, \dots$, consider q -exchangeable measure on all sequences of length n with precisely $k(n)$ 1's and $(n-k(n))$ 0's.

We need to find limits of such measures as $n \rightarrow \infty$

- $k(n)=0$. Then the only sequence is $(0, \dots, 0)$ and we converge to 0-measure of Theorem 1 as $n \rightarrow \infty$

- $k(n)=1$. Then we deal with measure with weights

$$\text{Prob} \left(\underbrace{0 \dots 0}_m \underbrace{1 0 \dots 0}_{n-m-1} \right) = \frac{1-q}{1-q^n} \cdot q^m, \quad 0 \leq m \leq n-1.$$

$\frac{1}{1+q+\dots+q^{n-1}}$

As $n \rightarrow \infty$ we converge to 1-measure of Theorem 1

- More generally, suppose that $k(n) \rightarrow k$ as $n \rightarrow \infty$. Then for large n we deal with measure on 0/1 sequences with 1's at positions $0 < y_1 < \dots < y_k \leq n$ and weight $\text{Prob}(y_1, \dots, y_k) = \frac{1}{Z} q^{y_1 + \dots + y_k}$ (*)

[Indeed, q -exchangeability of such measure is clear]

Here Z is a normalization constant, making total mass = 1. [In fact, $Z = q^{\binom{n}{k}}$ q -binomial coefficient.]

See q -Binomial theorem at https://en.wikipedia.org/wiki/Gaussian_binomial_coefficient

$$q^{y_1 + \dots + y_k} = q^{k \cdot y_1} \cdot q^{\binom{k-1}{1}(y_2 - y_1)} \cdot \dots \cdot q^{\binom{k-1}{k-1}(y_k - y_{k-1})} \quad (**)$$

This gives k -measure of Theorem 1 as $n \rightarrow \infty$

- The final case is $k(n) \rightarrow \infty$ as $n \rightarrow \infty$. We would like to show that in the limit we get ∞ -measure of Theorem 1.

Quitting some details, this can be seen either from (*):
 if, say, $y_1 > 1$, the $y_2 > 2, y_3 > 3, \dots$ and (*) gets multiplied
 at least by $q^{k(n)}$ [which is very small] compared to the
 minimal configuration $y_1 = 1, y_2 = 2, \dots$. Or from (**):
 $q^{k(n) \cdot y_1}$ factor makes $y_1 = 1$, $q^{(k(n)-1)(y_2-2)}$ factor makes
 $y_2 - y_1 = 1$ as $n \rightarrow \infty$, etc. □

Exercise 2: Check (directly from definition) that
 the measures of Theorem 1 are, indeed, q -exchangeable

Remark 2: For $q > 1$, the answer in an analogue
 of Theorem 1 is the same, but with roles of
 0's and 1's interchanged and q replaced by q^{-1}

There is also an analogue of Theorem 2 for $q > 1$.

Interpretation; infinite-dimensional Grassmannian over a finite field.

\mathbb{F}_q - finite field with q elements

V_n - n -dimensional vector space: (x_1, \dots, x_n) [$x_i \in \mathbb{F}_q$]

$V_\infty = \bigcup_{n=1}^{\infty} V_n$ - sequences (x_1, x_2, \dots) , such that $x_i = 0$ for large enough i .

[Such infinite-dimensional spaces with countable basis are rarely studied over \mathbb{R} , but there is no L_p over \mathbb{F}_q and we have no choice]

$Gr(V_\infty)$ - all (linear) subspaces in V_∞
Grassmannian

$GL(n, \mathbb{F}_q)$ - all invertible $n \times n$ matrices with elements from \mathbb{F}_q

$GL(\infty, \mathbb{F}_q) = \bigcup_{n=1}^{\infty} GL(n, \mathbb{F}_q)$ - infinite matrices of the form $\begin{pmatrix} \text{invertible} & 0 \\ 0 & 1 \dots \end{pmatrix}$

$GL(n, \mathbb{F}_q)$ acts in V_n . $GL(\infty, \mathbb{F}_q)$ acts in V_∞ and in $Gr(V_\infty)$

Random $GL(\infty, \mathbb{F}_q)$ -invariant subspace $X \in Gr(V_\infty)$?

Theorem 3: $GL(\infty, \mathbb{F}_q)$ -invariant probability measures on $Gr(V_\infty)$ are in bijection with q -exchangeable laws of sequences of 0/1. The correspondence is given by projection $\hat{\pi}: Gr(V_\infty) \rightarrow \{0, 1\}^\infty$

$$\hat{\pi}(X) = (\underbrace{\dim(X \cap V_1)}_{\text{dimension of space}}, \dim(X \cap V_2) - \dim(X \cap V_1), \dots, \dim(X \cap V_n) - \dim(X \cap V_{n-1}), \dots)$$

We do not give a proof — see Section 5 of Gnedin-Olshanski.

Idea: $X \leftrightarrow$ all $X \cap V_n$. $GL(n, \mathbb{F}_q)$ -orbit of $X \cap V_n$ is fixed by its dimension

Corollary: Ergodic random $GL(\infty, \mathbb{F}_q)$ -invariant subspaces X are:

- $X = \{0\}$ — subspace consisting of one element $(0, 0, \dots)$
- Random subspace of finite codimension $k=0, 1, 2, \dots$. In particular, for $k=0$, we have $X = V_\infty$.

Additional literature:

More general sequences than just 0/1 →

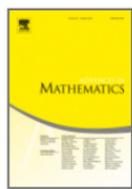
q-exchangeability via quasi-invariance

Alexander Gnedin, Grigori Olshanski

Ann. Probab. 38(6): 2103-2135 (November 2010). DOI: 10.1214/10-AOP536



Advances in Mathematics
Volume 254, 20 March 2014, Pages 331-395



← Asymptotic representation theory of $GL(n, \mathbb{F}_q)$

Finite traces and representations of the group of infinite matrices over a finite field

Vadim Gorin ^{a, b}, Sergei Kerov ¹, Anatoly Vershik ^{c, d}

← and more complicated measures on subspaces over \mathbb{F}_q

Random matrices over \mathbb{F}_q ↓



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Mathematics > Representation Theory

[Submitted on 3 Feb 2021]

Infinite-dimensional groups over finite fields and Hall-Littlewood symmetric functions

Cesar Cuenca, Grigori Olshanski

Random matrix theory over finite fields

Author: Jason Fulman

Journal: Bull. Amer. Math. Soc. 39 (2002), 51-85

Recall that at $q=1$: Pascal graph \hookrightarrow Gelfand-Tsetlin graph.

This embedding lifts to the q -level

Definition: q -Gelfand-Tsetlin graph q -GT has levels formed by signatures $\lambda_1 \geq \lambda_2 \geq \dots \geq \lambda_N$ and cotransition probabilities

$$L_{N+1 \rightarrow N}(\lambda \rightarrow \mu) = \begin{cases} (1-q) \dots (1-q^N) \cdot q^{\sum_{i=1}^N (\mu_i + N - i)} \cdot \frac{\prod_{i < j} (q^{\mu_i + N - i} - q^{\mu_j + N - i})}{\prod_{i < j} (q^{\lambda_j + N + 1 - j} - q^{\lambda_i + N + 1 - i})}, & \text{if } \mu \prec \lambda \\ 0, & \text{otherwise.} \end{cases}$$

interlace, i.e.
 $\lambda_1 \geq \mu_1 \geq \lambda_2 \geq \dots \geq \mu_N \geq \lambda_{N+1}$

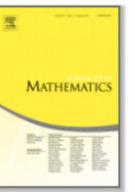
- q -Pascal is embedded a subgraph of signatures $(1 \geq 1 \geq \dots \geq 1 \geq 0 \dots \geq 0)$

[Exercise 1 on Slide 4 leads to that]

- At $q=1$ back to GT of Week 8

- This is also reminiscent of Week 10, Slide 10, Proposition 2

- The normalization $\sum_{\mu} L_{N+1 \rightarrow N}(\lambda \rightarrow \mu) = 1$ can be checked using Schur polynomials



The q -Gelfand-Tsetlin graph, Gibbs measures and q -Toeplitz matrices

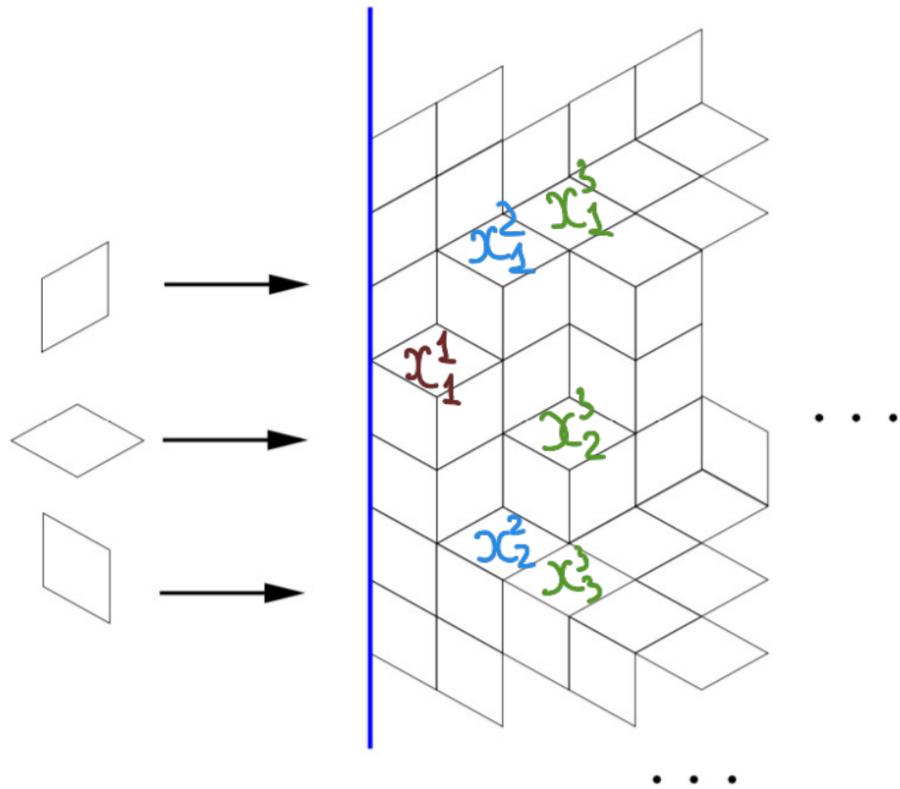
Vadim Gorin

Motivations for q -GT:

see Week 8, Slide 10

1. Combinatorial/probabilistic.

Recall that paths in GT (or q -GT) are in bijection with **lozenge tilings**, which can be also viewed as **two-dimensional stepped surfaces** in 3d-space.



Central measures in GT

Conditionally uniform random tilings

Multiplying links $L_{N+1 \rightarrow N} \cdot L_{N \rightarrow N-1} \cdots L_{2 \rightarrow 1}$, we get

Central measures in q -GT

Tilings with weight $q^{\sum_{k=1}^N \sum_{i=1}^k x_i^k}$ \sim $q^{\text{Volume under the surface}}$

This is the simplest possible non-uniform weight on tilings

2. Representation theoretic. Recall that for GT
 $\text{Dim}_N(\lambda)$ = dimension of irreducible representation of $U(N)$
 indexed by signature ("highest weight") $\lambda_1 \geq \lambda_2 \geq \dots \geq \lambda_N$

Similarly, for q -GT, setting $\text{Dim}_N^q(\lambda) = \sum_{\text{paths } \phi \rightarrow \lambda} q^{\text{Volume}}$

we have $\text{Dim}_N^q(\lambda)$ = quantum dimension of representation
 of quantized universal enveloping algebra ("quantum group")
 $U_q(\mathfrak{gl}_N)$ indexed by signature $\lambda_1 \geq \dots \geq \lambda_N$

Such representation has the same dimension as vector space,
 as the $U(N)$ representation. However, there is a distinguished
 element q^{β} in quantum group and by many reasons it is more
 natural to consider its trace in representation (= quantum dimension).

Similarly, there are quantum traces: $\text{Trace} \left(\overset{\text{usual trace}}{\text{T}} \left(\overset{\text{representation}}{A} q^{\beta} \right) \right)$
 where β is an element of quantum group.

Theorem 4: Assume $0 < q < 1$. Extreme coherent systems in q -GT are parameterized by infinite sequences of integers: $J_1 \leq J_2 \leq J_3 \leq \dots \in \mathbb{N}$

Let $\lambda^{(1)} \prec \lambda^{(2)} \prec \lambda^{(3)} \prec \dots$ be a path distributed according to $(q-)$ central measure corresponding to such sequence. Then

$$\lim_{N \rightarrow \infty} \lambda_{N+1-i}^{(N)} = J_i \quad [\text{No normalizations!}]$$

Suppose that $J_1 \geq 0$. Then $\lambda_1^{(1)} \geq 0$ almost surely and

$$\sum_{m=0}^{\infty} \text{Prob}(\lambda_1^{(1)} = m) \underbrace{(1-x)(1-qx) \dots (1-q^{m-1}x)}_{q\text{-interpolation polynomial}} = \prod_{i \geq 1} \frac{1 - q^{i-1}x}{1 - q^{J_i + i - 1}x}$$

- The structure of q -GT is invariant under simultaneous shifts of all coordinates $\lambda_i^k \mapsto \lambda_i^k + c$. Same shifts apply to $J_1 \leq J_2 \leq \dots$
- The law of $(\lambda_1^k, \dots, \lambda_k^k)$ can be similarly described in terms of q -Interpolation Schur polynomials

Four proofs of Theorem 4 in the literature



Advances in Mathematics
Volume 229, Issue 1, 15 January 2012, Pages 201-266



The q -Gelfand–Tsetlin graph, Gibbs measures and q -Toeplitz matrices

Vadim Gorin

Moscow Mathematical Journal

Volume 14, Issue 1, January–March 2014 pp. 121–160.

The Boundary of the Gelfand–Tsetlin Graph: New Proof of Borodin–Olshanski’s Formula, and its q -Analogue

Authors: Leonid Petrov

The Annals of Probability
2015, Vol. 43, No. 6, 3052–3132
DOI: 10.1214/14-AOP955
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ASYMPTOTICS OF SYMMETRIC POLYNOMIALS WITH
APPLICATIONS TO STATISTICAL MECHANICS AND
REPRESENTATION THEORY

BY VADIM GORIN¹ AND GRETA PANOVA²



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Volume 270, Issue 1, 1 January 2016, Pages 375–418



A quantization of the harmonic
analysis on the infinite-dimensional
unitary group

Vadim Gorin^{a, b}, Grigori Olshanski^{b, c, d}

We will only sketch some ideas.

As usual, we compute the Martin boundary first and start from $\lambda_1 \geq \dots \geq \lambda_N$. We act on it with $(q-)$ links to get a (random) signature $\mu_1 \geq \dots \geq \mu_k$. We aim to send $N \rightarrow \infty$ keeping k fixed. Note that due to interlacing, $\mu_k \geq \lambda_N, \mu_{k-1} \geq \lambda_{N-1}, \dots, \mu_1 \geq \lambda_{N-k+1}$ almost surely.

This was also true at $q=1$ and in β -spectra graph. However, the new feature of $0 < q < 1$ setting is that

$\text{Prob}(\mu_k = \lambda_N; \mu_{k-1} = \lambda_{N-1}; \dots; \mu_1 = \lambda_{N-k+1}) > C > 0$ with C staying away from 0 as $N \rightarrow \infty$.

This can be proven purely probabilistically by analyzing the cotransition probabilities, and it readily implies the necessity of $\lambda(N)_{N-i+1} \rightarrow \downarrow_i$, $i=1,2,\dots$
 [as otherwise a part of the measure is escaping to $\pm\infty$]

However, this approach does not give any formulas for the limiting coherent systems (Martin boundary)

The latter can be obtained by using Schur functions [see week 6, slide 5]

Proposition 1: For any $\lambda = (\lambda_1 \geq \dots \geq \lambda_{N+1})$ we have

$$\frac{S_\lambda(x_1, \dots, x_N, q^{-N})}{S_\lambda(1, q^{-1}, \dots, q^{-N})} = \sum_{\mu = (\mu_1 \geq \dots \geq \mu_N)} \overset{\text{in } q\text{-GT}}{\downarrow} L_{(N+1) \rightarrow N}(\lambda \rightarrow \mu) \frac{S_\mu(x_1, \dots, x_N)}{S_\mu(1, q^{-1}, \dots, q^{1-N})}$$

Thus, cotransition probabilities are encoded through branching rules for Schur polynomials upon restrictions to smaller # of variables.

We need auxiliary lemmas:

Lemma 1: $S_\lambda(1, q, \dots, q^{N-1}) = \prod_{i < j} \frac{q^{\lambda_i + N - i} - q^{\lambda_j + N - j}}{q^{N-i} - q^{N-j}}$

Proof. We use **Vandermonde determinant**: $\det [x_i^{N-j}]_{i,j=1}^N = \prod_{i < j} (x_i - x_j)$

Then $S_\lambda(1, \dots, q^{N-1}) = \frac{\det [q^{(N-j)(\lambda_i + N - i)}]}{\prod_{i < j} (q^{N-i} - q^{N-j})} = \prod_{i < j} \frac{q^{\lambda_i + N - i} - q^{\lambda_j + N - j}}{q^{N-i} - q^{N-j}}$ \square

Lemma 2. $S_\lambda(ax_1, \dots, ax_n) = a^{\lambda_1 + \dots + \lambda_n} S_\lambda(x_1, \dots, x_n)$
[Immediate from definition]

Lemma 3. $S_\lambda(x_1, \dots, x_n, t) = \sum_{\mu \prec \lambda} t^{\sum \lambda_i - \sum \mu_i} S_\mu(x_1, \dots, x_n)$
interlace $\rightarrow \mu \prec \lambda$

For the proof of **Lemma 3**, repeat the argument of **Week 11**, slides 16-17 or see **MATH 740**, Lecture 4, slides 3-5

Proof of Proposition 1: The last lemma implies

$$\frac{S_\lambda(x_1, \dots, x_N, q^{-N})}{S_\lambda(1, q^{-1}, \dots, q^{-N})} = \sum_{\mu \prec \lambda} q^{-N(\sum_{i=1}^{N+1} \lambda_i - \sum_{i=1}^N \mu_i)} \cdot \frac{S_\mu(1, \dots, q^{1-N})}{S_\lambda(1, \dots, q^{-N})} \cdot \frac{S_\mu(x_1, \dots, x_N)}{S_\mu(1, \dots, q^{1-N})}$$

It remains to identify $L_{(N+1) \rightarrow N}(\lambda \rightarrow \mu) = q^{-N(\sum_{i=1}^{N+1} \lambda_i - \sum_{i=1}^N \mu_i)} \cdot \frac{S_\mu(1, \dots, q^{1-N})}{S_\lambda(1, \dots, q^{-N})}$ using Lemmas 1 and 2. \square

Corollary 1: $\sum_{\mu} L_{(N+1) \rightarrow N}(\lambda \rightarrow \mu) = 1.$

Proof: Set $(x_1, \dots, x_N) = (1, q^{-1}, \dots, q^{1-N})$ in Proposition \square

Corollary 2: Fix $\lambda = (\lambda_1, \dots, \lambda_N)$ and let $\lambda^{(1)} \prec \lambda^{(2)} \prec \dots \prec \lambda^{(n)} = \lambda$ be a q^{vol} -weighted random path. [Trajectory of chain $L_{N \rightarrow N-1}; L_{N-1 \rightarrow N-2}; \dots$]

Then $\sum_{m \in \mathbb{N}} \text{Prob}(\lambda^{(n)} = m) x^m = \frac{S_\lambda(x, q^{-1}, \dots, q^{1-N})}{S_\lambda(1, q^{-1}, \dots, q^{1-N})}$

Proof. We iterate Proposition 1 ($N-1$) times. \square

Remark 2: Setting $q=1$ in Corollary 2 we get Week 8, Slide 8, Th. 4

Remark 3: There is an immediate analogue of Corollary 2 for the generating function of distribution of $(\lambda_1^{(n)}, \dots, \lambda_n^{(n)})$
[Write it down!]

Conclusion: In order to find the Martin boundary of q -GT, we need to find limits of $\frac{S_\lambda(x, q^{-1}, \dots, q^{1-N})}{S_\lambda(1, \dots, q^{1-N})}$ as $N \rightarrow \infty$ and show that the last coordinates of λ should stabilize to $\lambda_1, \lambda_2, \dots$

Proposition 2:
$$\frac{S_\lambda(x, q^{-1}, \dots, q^{1-N})}{S_\lambda(1, \dots, q^{1-N})} = \left[\prod_{i=1}^{N-1} \frac{1 - q^i}{1 - x q^i} \right] \cdot \frac{\ln q}{2\pi i} \cdot \oint_{\{\lambda_i + N - i\}} \frac{x^z dz}{\prod_{i=1}^N (1 - q^{-z} q^{\lambda_i + N - i})}$$

where the integration goes over a positively \circlearrowright oriented contour enclosing the poles at $\lambda_i + N - i$, $i = 1, 2, \dots, N$.

Proof of Proposition 2: $S_\lambda(x_1, \dots, x_N) \cdot \prod_{i < j} (x_i - x_j) = \det [x_i^{\lambda_j + N - j}]$
 $= \sum_{k=1}^N (-1)^{k-1} x_1^{\lambda_k + N - k} \cdot \det [x_i^{\lambda_j + N - j}]_{i \geq 2, j \neq k}$

We can plug $x_i = q^{1-i}$, $i = 2, 3, \dots, N$, and $x_1 = x$.
 Determinant is again Vandermonde, as in Lemma 1, and is, therefore,
 explicitly evaluated. Hence, $\frac{S_\lambda(x, \dots, q^{1-N})}{S_\lambda(1, \dots, q^{1-N})}$ is an explicit sum of N terms.

These terms match the expansion of the contour integral as a sum
 of N residues at poles $z = \lambda_j + N - j$. 

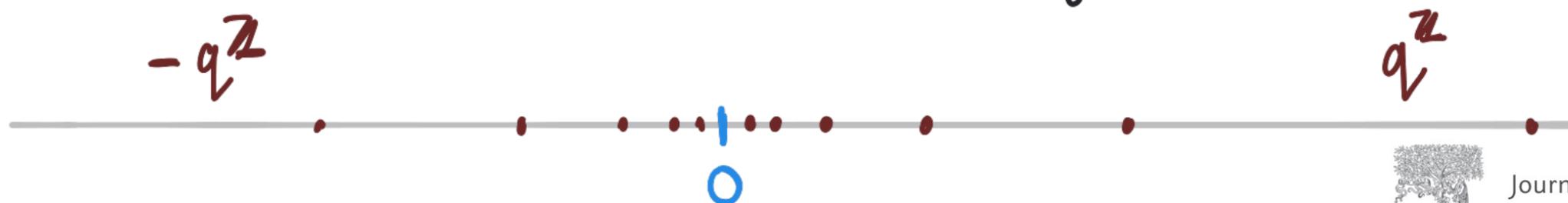
Remark 4: Sending $q \rightarrow 1$, we get formulas of Week 8, slide 9

Sending $N \rightarrow \infty$ in the result of Proposition 2 is not hard:
 The integrand converges towards $\frac{x^z}{\prod_{i=1}^{\infty} (1 - q^{-z} q^{j_i + i - 1})}$, and
 the prefactor becomes $\prod_{i=1}^{\infty} \frac{1 - q^i}{1 - x q^i}$. [The contour needs a deformation to make sense as $n \rightarrow \infty$]
 Resummations eventually lead to the result of Theorem 4.

We do not provide further details on **Theorem 4**, but let us mention that it has a remarkable generalization.

If we interpret $(\lambda_1, \dots, \lambda_N)$ as N -particle configuration $\{q^{\lambda_i + N - i}\}_{i=1}^N$ on the lattice $q^{\mathbb{Z}}$, then the link of **Slide 13** takes the form $L_{(n+1) \rightarrow n}(X \rightarrow Y) =$
 $= (1-q) \dots (1-q^N) \prod_{i=1}^N |y_i| \cdot \frac{\prod_{i < j} (y_j - y_i)}{\prod_{i < j} (x_i - x_j)}$ [Y interlaces with X]

One can extend the same formula to configurations on **two-sided lattice**.



- Analogue of **Theorem 4**
- Important coherent systems, which do not exist on q -GT
- Representation-theoretic meaning of doubling is **unknown**

